A LIOUVILLE THEOREM FOR F-HARMONIC MAPS WITH MODERATE DIVERGENT ENERGY

HO CHOR YIN

Abstract. In this paper, we obtain a Liouville theorem of F-harmonic maps between complete Riemannian manifolds with moderate divergent F-energy. We assume that F is a concave function and satisfies a differential inequality. We employ Ara's F-stress-energy tensor and the Hessian comparison theorem to prove the main result.

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1. INTRODUCTION

Harmonic maps [5,11] are critical points of the energy functional defined on the space of smooth maps between Riemannian manifolds. In [1], the author introduced the theory of F-harmonic maps to unify various types of harmonic map, e.g. p-harmonic maps and exponentially harmonic maps. Liouville type properties of harmonic maps were studied by several authors.(cf. [3,8,9] and the reference therein). In this paper, we study Liouville type properties of Fharmonic maps with moderate divergent F-energy (see Definition 2.5). That is, we study conditions for which a F-harmonic map with moderate divergent F-energy u between two Riemannian manifolds (M, g) and (N, h) is a constant map. In [11, p. 46], the author obtained the following

THEOREM 1.1 ([11]). Let M be a Cartan-Hadamard manifold of dimension m whose sectional varies in a small range, and u a harmonic map from M into any Riemannian manifold N with moderate divergent energy. If the dimension m of the domain manifold is greater than 2, then u has to be constant.

Recall that a Cartan-Hadamard manifold is a complete simply-connected, non-positive sectional curvature Riemannian manifold. We want to know whether for F-harmonic map, we can have similar Liouville type properties. In [4], the authors performed deep analysis about the growth rate of F-energy and used asymptotic assumption of the map at infinity to obtain a series of

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Liouville theorems under various conditions e.g. pinching of radial curvature of the domain manifold. In [10], the author obtained some Liouville theorems under finiteness F-energy condition. To the best of our knowledge, Liouville theroem for F-harmonic maps with the moderate energy divergent condition (In the sense of Xin's works) have not appeared in the literature. A key ingredient in Xin's proof is the conservation law of u in terms of the the vanishing of the divergence of the stress energy tensor of harmonic map. Therefore, we need to find a conservation law for F-harmonic map with a suitably defined F-stress energy tensor. It is known that the Ara's stress-energy tensor (See (7)) for F-harmonic map is conserved [1] and we shall use this tensor in our study. Our main theorem is the following

THEOREM 1.2. Let $(M^m, g), m > 2$ be a complete, simply connected with nonnegative sectional curvature Riemannian manifold which has a pole and (N^n, h) be any Riemannian manifold. Let $u : (M^m, g) \to (N^n, h)$ be Fharmonic map where $F : (0, \infty) \to [0, \infty)$ is a C^2 function such that F(0) = 0, F' > 0 on $(0, \infty)$ and F is a concave function satisfying the differential inequality,

(1)
$$tF'(t)/F(t) \le C$$

where C is a positive constant. Assume that the sectional curvature of (M^m, g) varies in a small range, more precisely, we assume that the radial curvature k of the domain manifold M satisfies $-c_1^2 \leq k \leq -c_2^2 < 0$, where c_1, c_2 are positive constants, and with moderate divergent F-energy. Then the map u has to be constant.

To replace the condition finiteness of F-energy with the condition of moderate divergent F-energy, we assume that the function F(t) is concave and satisfies a differential inequality. To prove the main result, we firstly obtain an identity about the divergence of the F-stress energy tensor and take integration to obtain Lemma 2.4; then we apply Hessian comparison theorem to get some key estimates. Throughout this paper, we will adopt Einstein summation convention: sum on repeated indices. This paper is organized as follows, in Section 2, we recall definition of F-harmonic maps, conservation law and stress energy tensor. In Section 3, we present the proof of the main results.

2. PRELIMINARIES

2.1. F-ENERGY AND F-HARMONIC MAPS

Let $F : \mathbb{R} \to \mathbb{R}$ be a C^{∞} function. For a smooth map u between Riemannian manifolds (M, g) and (N, h), the F-energy $E_F(u)$ is defined by

(2)
$$E_F(u) = \int_M F\left(\frac{|\mathrm{d}u|^2}{2}\right) * 1.$$

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While let us recall the energy of E(u) is defined by

(3)
$$E(u) = \int_M \left(\frac{|\mathrm{d}u|^2}{2}\right) * 1 = \int_M \left(\frac{1}{2}h_{\alpha\beta}\frac{\partial u^\alpha}{\partial x^i}\frac{\partial u^\beta}{\partial x^j}g^{ij}\right) * 1$$

Here, the volume element

(4)
$$*1 = \sqrt{|g|} \mathrm{d}x^1 \wedge \cdots \wedge \mathrm{d}x^m.$$

A smooth map $u: (M,g) \to (N,h)$ is said to be *F*-harmonic if $u|_{\Omega}$ is a critical point of the *F*-energy $E_f(u)$ for every compact domain $\Omega \in M$.

2.2. CONSERVATION LAW FOR F-HARMONIC MAP

Let $u : (M,g) \to (N,h)$ be a smooth map. Let \mathcal{E} and u^*h denote the energy density and the first fundamental form, respectively. Let $\mathcal{E}(u) := \frac{|du|^2}{2}, \mathcal{E}_F(u) := F(\frac{|du|^2}{2})$. For harmonic maps, the stress-energy(SE) tensor S_u is defined by

(5)
$$S_u = \mathcal{E}(u)g - u^*h,$$

which is a symmetric 2-tensor. Here, $\mathcal{E}(u)$ denotes the energy density function. A detailed treatment on stress-energy tensor of harmonic map can be found in [2,5]. For harmonic maps between manifolds, the basic relation between the stress-energy tensor and harmonic maps is the following [5]

PROPOSITION 2.1. div $S_u = -\langle \tau_u, \mathrm{d}u \rangle$.

DEFINITION 2.2. Let $u : (M^m, g) \to (N^n, h)$ be any smooth map. If u satisfies div $S_u \equiv 0$, then we say that the map u satisfies the conservation law.

The Euler-Lagrange equation of F-harmonic map gives us,

(6)
$$\tau_u^F = F'(\mathcal{E}(u))\tau_u + u_*\left(\operatorname{grad}(F'(\mathcal{E}(u)))\right).$$

and we have the following F-stress-energy tensor [1],

(7)
$$S_u^F = F(\mathcal{E}(u))g - F'(\mathcal{E}(u))u_*h.$$

Then, it holds that

(8)
$$\operatorname{div} S_u^F = -\langle \tau_u^F, \mathrm{d}u \rangle.$$

REMARK 2.3. We say that the F-stress energy tensor is conserved or satisfies a conservation law since it satisfies the equation(8).

LEMMA 2.4. Assume that $D \subseteq M$ is a geodesic ball and its boundary ∂D is a geodesic sphere, and u is F-harmonic. Choose a local orthonormal frame field of M along ∂D such that $f_1, \ldots, f_{m-1} \in \Gamma(T\partial D)$, f_m is the normal vector of ∂D , $\mathbf{n} := f_m$, we have

(9)
$$\int_{\partial D} \left(F(\mathcal{E}(u)) \langle X, \mathbf{n} \rangle - F'(\mathcal{E}(u)) \langle u_* X, u_* \mathbf{n} \rangle) \right) * 1 = \int_D \langle S_u^F, \nabla X \rangle * 1,$$

where

(10)
$$\langle S_u^F, \nabla X \rangle = F(\mathcal{E}(u)) \operatorname{div} X - F'(\mathcal{E}(u)) \langle u_* e_i, u_* e_j \rangle_h \langle \nabla_{e_i} X, e_j \rangle_g.$$

Proof. Let $\{e_i\}$ be an orthonormal frame near a point $p \in M$ such that $\nabla_{e_i} e_j|_p = 0$. Let $X \in \Gamma(TM)$ be any vector field on M, we have $\operatorname{div}(F(\mathcal{E}(u)X) = X(F(\mathcal{E}(u))) + F(\mathcal{E}(u))\operatorname{div} X$

(11)

$$div(F(\mathcal{E}(u)X) = X(F(\mathcal{E}(u))) + F(\mathcal{E}(u))divX$$

$$= F'(\mathcal{E}(u))X(\mathcal{E}(u)) + F(\mathcal{E}(u))divX$$

$$= F'(\mathcal{E}(u))X(\mathcal{E}(u)) + \langle \nabla X, F(\mathcal{E}(u))g \rangle,$$

where we use the notation

(12)
$$\nabla X(e_i, e_j) := \langle \nabla_{e_i} X, e_j \rangle.$$

In the following, we will use $du(e_i) = u_*e_i$ interchangeably for simplicity. We compute the first term in (11) as follows,

$$\begin{split} F'(\mathcal{E}(u))X(\mathcal{E}(u)) &= F'(\mathcal{E}(u))\frac{1}{2}\nabla_X \langle u_*e_i, u_*e_i \rangle \\ &= F'(\mathcal{E}(u)) \langle \nabla_X u_*e_i, u_*e_i \rangle \\ &= F'(\mathcal{E}(u)) \langle (\nabla_x du)e_i, u_*e_i \rangle \\ &= F'(\mathcal{E}(u)) \langle (\nabla_{e_i} du)X, u_*e_i \rangle \\ &= \langle (\nabla_{e_i} du)X, F'(\mathcal{E}(u))u_*e_i \rangle - \langle du(\nabla_{e_i}X), F'(\mathcal{E}(u))u_*e_i \rangle \\ &= \nabla_{e_i} \langle du(X), F'(\mathcal{E}(u))u_*e_i \rangle - \langle du(X), \nabla_{e_i}(F'(\mathcal{E}(u))u_*e_i) \rangle \\ &- \langle du(\nabla_{e_i}X), F'(\mathcal{E}(u))u_*e_i \rangle \\ &= \nabla_{e_i} \langle du(X), F'(\mathcal{E}(u))u_*e_i \rangle \\ &= div \left(F'(\mathcal{E}(u)) \langle du(\nabla_{e_i}X), u_*e_i \rangle \\ \\ &= div \left(F'(\mathcal{E}(u)) \langle u_*X, u_*e_i \rangle e_i\right) - \langle u_*X, \tau_u^F \rangle \\ &= div \left(F'(\mathcal{E}(u)) \langle u_*X, u_*e_i \rangle e_i\right) - \langle u_*X, \tau_u^F \rangle \\ &= div \left(F'(\mathcal{E}(u)) \langle u_*X, u_*e_i \rangle e_i\right) - \langle u_*X, \tau_u^F \rangle \\ &= div \left(F'(\mathcal{E}(u)) \langle u_*X, u_*e_i \rangle e_i\right) - \langle u_*X, \tau_u^F \rangle \\ &= div \left(F'(\mathcal{E}(u)) \langle v_X, u^*h \rangle, \end{split}$$

substituting this into (11), we have

(13)

$$div(F(\mathcal{E}(u)X) = div(F'(\mathcal{E}(u))\langle u_*X, u_*e_i\rangle e_i) - \langle u_*X, \tau_u^F \rangle - F'(\mathcal{E}(u))\langle \nabla X, u^*h \rangle + \langle \nabla X, F(\mathcal{E}(u))g \rangle,$$

$$= div(F'(\mathcal{E}(u))\langle u_*X, u_*e_i\rangle e_i) - \langle u_*X, \tau_u^F \rangle + \langle S_u^F, \nabla X \rangle.$$

As $D \subseteq M$ is a geodesic ball and its boundary is a geodesic sphere, and u is a F-harmonic map and $\{e_i\}$ be an orthonormal frame near a point $p \in \partial D$ such that $e_1, \ldots, e_{m-1} \in \Gamma(T \partial D)$, e_m is the normal vector of ∂D , $\mathbf{n} := e_m$ and $\nabla_{e_i} e_j |_p = 0$, taking integration on D for (13), we obtain

$$\int_{\partial D} \left(F(\mathcal{E}(u)) \langle X, \mathbf{n} \rangle - F'(\mathcal{E}(u)) \langle u_* X, u_* \mathbf{n} \rangle) \right) * 1 = \int_D \langle S_u^F, \nabla X \rangle * 1.$$

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We modify the definition of the condition of moderate divergent energy in [11] for F-energy to obtain the following

DEFINITION 2.5. If there is a positive function $\eta(r)$ and $R_0 > 0$ such that

$$\int_{R_0}^{\infty} \frac{\mathrm{d}r}{r\eta(r)} = \infty,$$

then we have

$$\lim_{R \to \infty} \int_{B_R(x_0)} \frac{F(\mathcal{E}(u))(x)}{\eta(r(x))} * 1 < \infty.$$

Then, we say that u is moderate F-energy divergent.

3. PROOF OF MAIN RESULTS

In this section, we will give the proof of Theorem 1.2.

Proof. Let $D = B_R(x_0)$ be a geodesic ball of radius R and centered at x_0 and its boundary is the geodesic sphere $\partial B_R(x_0)$. It is clear that the square of the distance function from x_0 in $B_R(x_0)$ is smooth. Let $\frac{\partial}{\partial r}$ be the unit radial vector field which is also the unit normal vector field \mathbf{n} to $\partial B_R(x_0)$. Take $X = r \frac{\partial}{\partial r}$ in (9). L.H.S of (9) becomes

$$\int_{\partial B_R(x_0)} \left(F(\mathcal{E}(u)) \langle X, \mathbf{n} \rangle - F'(\mathcal{E}(u)) \langle u_* X, u_* \mathbf{n} \rangle \right) \right) * 1$$

=
$$\int_{\partial B_R(x_0)} RF(\mathcal{E}(u)) * 1 - \int_{\partial B_R(x_0)} F'(\mathcal{E}(u)) R \langle u_* \frac{\partial}{\partial r}, u_* \frac{\partial}{\partial r} \rangle) * 1$$

$$\leq R \int_{\partial B_R(x_0)} F(\mathcal{E}(u)) * 1.$$

We want to obtain a lower bound of the R.H.S. of (9), to this end, we compute the following first,

(14)
$$\nabla_{\frac{\partial}{\partial r}} X = \nabla_{\frac{\partial}{\partial r}} (r \frac{\partial}{\partial r}) = \frac{\partial}{\partial r},$$
$$\nabla_{e_s} X = \nabla_{e_s} (r \frac{\partial}{\partial r}) = r \nabla_{e_s} \frac{\partial}{\partial r}$$

(15)
$$= r \operatorname{Hess}(r)(e_s, e_t)e_t,$$

(16)
$$\operatorname{div} X = 1 + r \operatorname{Hess}(r)(e_s, e_s)$$

where $\{e_{\alpha}\} := \{e_s, \frac{\partial}{\partial r}\}$ is a local orthonormal frame near x_0 on $B_R(x_0)$. Using (15),(16), we have

$$\langle u_* e_{\alpha}, u_* e_{\beta} \rangle \langle \nabla_{e_{\alpha}} X, e_{\beta} \rangle$$

$$= \langle u_* e_s, u_* e_t \rangle \langle \nabla_{e_s} X, e_t \rangle + \langle u_* \frac{\partial}{\partial r}, u_* e_t \rangle \langle \nabla_{e_s} X, e_t \rangle$$

$$+ \langle u_* e_s, u_* \frac{\partial}{\partial r} \rangle \langle \nabla_{e_s} X, \frac{\partial}{\partial r} \rangle + \langle u_* \frac{\partial}{\partial r}, u_* \frac{\partial}{\partial r} \rangle \langle \nabla_{\frac{\partial}{\partial r}} X, \frac{\partial}{\partial r} \rangle$$

$$= \langle u_* e_s, u_* e_t \rangle \langle \nabla_{e_s} X, e_t \rangle + \langle u_* \frac{\partial}{\partial r}, u_* \frac{\partial}{\partial r} \rangle$$

$$= r \operatorname{Hess}(r)(e_s, e_t) \langle u_* e_s, u_* e_t \rangle + \langle u_* \frac{\partial}{\partial r}, u_* \frac{\partial}{\partial r} \rangle.$$

Using (10),(17), we obtain

(18)

$$\begin{aligned} \langle S_u^F, \nabla X \rangle &= F(\mathcal{E}(u)) \operatorname{div} X - F'(\mathcal{E}(u)) \langle u_* e_\alpha, u_* e_\beta \rangle_h \langle \nabla_{e_\alpha} X, e_\beta \rangle_g \\ &= F(\mathcal{E}(u)) (1 + r \operatorname{Hess}(r)(e_s, e_s)) \\ &- F'(\mathcal{E}(u)) \left(r \operatorname{Hess}(r)(e_s, e_t) \langle u_* e_s, u_* e_t \rangle + \langle u_* \frac{\partial}{\partial r}, u_* \frac{\partial}{\partial r} \rangle \right).
\end{aligned}$$

We consider the case that when the radial curvature k of the domain manifold M satisfies $-c_1^2 \leq k \leq -c_2^2 < 0$, where c_1, c_2 are positive constants. Using the Hessian comparison theorem [7], we see that (18) becomes

$$\begin{split} \langle S_u^F, \nabla X \rangle &\geq F(\mathcal{E}(u))[1 + (m-1)(c_2 r) \coth(c_2 r)] \\ &- F'(\mathcal{E}(u)) \left| u_* \frac{\partial}{\partial r} \right|^2 - F'(\mathcal{E}(u))(c_1 r) \coth(c_1 r) \langle u_* e_s, u_* e_s \rangle \\ &= F\left(\frac{1}{2} \langle u_* e_s, u_* e_s \rangle + \frac{1}{2} \langle u_* \frac{\partial}{\partial r}, u_* \frac{\partial}{\partial r} \rangle \right) [1 + (m-1)(c_2 r) \coth(c_2 r)] \\ &- F'(\mathcal{E}(u)) \left| u_* \frac{\partial}{\partial r} \right|^2 - F'(\mathcal{E}(u))(c_1 r) \coth(c_1 r) \langle u_* e_s, u_* e_s \rangle \\ &\geq \left\{ F\left(\frac{1}{2} \langle u_* e_s, u_* e_s \rangle \right) + F\left(\frac{1}{2} \langle u_* \frac{\partial}{\partial r}, u_* \frac{\partial}{\partial r} \rangle \right) \right\} \\ &[1 + (m-1)(c_2 r) \coth(c_2 r)] \\ &- F'(\mathcal{E}(u)) \left| u_* \frac{\partial}{\partial r} \right|^2 - F'(\mathcal{E}(u))(c_1 r) \coth(c_1 r) \langle u_* e_s, u_* e_s \rangle \\ &\geq [F(\frac{1}{2} \langle u_* e_s, u_* e_s \rangle) + F(\frac{1}{2} \left| u_* \frac{\partial}{\partial r} \right|^2)] \\ &[1 + (m-1)(c_2 r) \coth(c_2 r)] \\ &- F'(\frac{1}{2} \left| u_* \frac{\partial}{\partial r} \right|^2) (\left| u_* \frac{\partial}{\partial r} \right|^2) \end{split}$$

$$-F'(\frac{1}{2}\langle u_*e_s, u_*e_s\rangle)(c_1r) \coth(c_1r)(\langle u_*e_s, u_*e_s\rangle)$$

$$\geq F(\frac{1}{2}\langle u_*e_s, u_*e_s\rangle) + r \coth(c_2r)[(m-1)c_2F(\frac{1}{2}\langle u_*e_s, u_*e_s\rangle)]$$

$$-\langle u_*e_s, u_*e_s\rangle c_1F'(\frac{1}{2}\langle u_*e_s, u_*e_s\rangle)] + F(\frac{1}{2}\left|u_*\frac{\partial}{\partial r}\right|^2)$$

$$-F'(\frac{1}{2}\left|u_*\frac{\partial}{\partial r}\right|^2)(\left|u_*\frac{\partial}{\partial r}\right|^2)$$

$$\geq B_1[F(\frac{1}{2}\langle u_*e_s, u_*e_s\rangle) + F(\frac{1}{2}\left|u_*\frac{\partial}{\partial r}\right|^2)]$$

where B_1 is a positive constant. Using (9) and (19), we obtain

(20)

$$\frac{R \int_{\partial B_R(x_0)} F(\mathcal{E}(u)) * 1}{\geq \int_{B_R(x_0)} B_1[F(\frac{1}{2} \langle u_* e_s, u_* e_s \rangle) + F(\frac{1}{2} \left| u_* \frac{\partial}{\partial r} \right|^2)] * 1.}$$

We will finish the proof by contradiction. If u is not a constant map then $F(\frac{1}{2}\langle u_*e_s, u_*e_s\rangle) + F(\frac{1}{2}|u_*\frac{\partial}{\partial r}|^2)$ is not identically equal to zero. There exists $R_0 > 0$, such that whenever $R > R_0$, we have

(21)
$$\int_{B_R(x_0)} \left[F(\frac{1}{2} \langle u_* e_s, u_* e_s \rangle) + F(\frac{1}{2} \left| u_* \frac{\partial}{\partial r} \right|^2) \right] * 1 \ge B_2,$$

where B_2 is a positive constant. Combining (20) and (21), we have

(22)
$$\int_{\partial B_R(x_0)} F(\mathcal{E}(u)) * 1 \ge \frac{B_1 B_2}{R}.$$

As we assume that u is moderate F-energy divergent, we recall that if there is a positive function $\eta(r)$ and $R_0 > 0$ such that

(23)
$$\int_{R_0}^{\infty} \frac{\mathrm{d}r}{r\eta(r)} = \infty,$$

then we have

(24)
$$\lim_{R \to \infty} \int_{B_R(x_0)} \frac{F(\mathcal{E}(u))(x)}{\eta(r(x))} * 1 < \infty.$$

Now, by using (24), (22) implies that

$$\lim_{R \to \infty} \int_{B_R(x_0)} \frac{F(\mathcal{E}(u))(x)}{\eta(r(x))} * 1 = \int_0^\infty \frac{\mathrm{d}r}{\eta(r)} \int_{B_R(x_0)} F(\mathcal{E}(u)) * 1$$
$$\geq B_1 B_2 \int_0^\infty \frac{\mathrm{d}r}{r\eta(r)}$$
$$\geq B_1 B_2 \int_{R_0}^\infty \frac{\mathrm{d}r}{r\eta(r)} = \infty.$$

This contradicts to the fact that u is moderate F-energy divergent. We conclude that u has to be a constant map.

COROLLARY 3.1. Assume that $(M^m, g), m > 2$ is a complete, simply connected with nonnegative sectional curvature Riemannian manifold which has a pole and (N^n, h) be any Riemannian manifold. Let $u : (M^m, g) \to (N^n, h)$ be F-harmonic map where $F : (0, \infty) \to [0, \infty)$ is a C^2 function such that F(0) = 0, F' < 0 Assume that the sectional curvature k of (M^m, g) satisfies: $-c_1^2 \leq k \leq -c_2^2 < 0$, where c_1, c_2 are positive constants., with moderate divergent energy. Then u has to be constant.

Proof. The setup is similar to that of the previous proof. As before, using the Hessian comparison theorem, we see that (18) becomes

$$\begin{aligned} \langle S_u^F, \nabla X \rangle &\geq F(\mathcal{E}(u))[1 + (m-1)(c_2 r) \coth(c_2 r)] \\ &- F'(\mathcal{E}(u)) \left| u_* \frac{\partial}{\partial r} \right|^2 - F'(\mathcal{E}(u))(c_1 r) \coth(c_1 r) \langle u_* e_s, u_* e_s \rangle \\ &\geq CF(\mathcal{E}(u)). \end{aligned}$$

The remaining steps are similar to the proof of Theorem 1.2 and we omit them. $\hfill \Box$

REFERENCES

- [1] M. Ara, Geometry of F-harmonic maps, Kodai Math. J., 22 (1999), 242–263.
- [2] P. Bair and J. Eells, A conservation law for harmonic maps, in Geometry Symposium Utrecht, 1980.
- [3] S. Y. Cheng, Liouville theorem for harmonic maps, in Proc. Sympos. Pure Math., Amer. Math. Soc., 36 (1980), 147–151.
- [4] Y. Dong, H. Lin and G. Yang, *Liouville theorems for F-harmonic maps*, Results Math., 69 (2016), 105-–127.
- [5] J. Eells and L. Lemaire, Selected Topics in Harmoinc Maps, in CBMS Regional Conference Series in Mathematics, Ed. 50, 1983.
- [6] J. Eells and J. Sampson, Harmonic maps of Riemannian manifolds, Amer. J. Math., 86 (1964), 109–164.
- [7] R. E. Greene and H. Wu, Function Theory on Manifolds Which Posses a Pole, Lecture Notes in Math., Vol. 699, Springer-Verlag, Berlin, 1979.
- [8] S. Hildebrandt, Liouville theorems for harmonic mappings, and an approach to Bernstein theorems, Annals of Mathematics Studies, Vol. 102, 1982, 107–131.
- [9] Z. Jin, Liouville theorems for harmonic maps, Invent. Math., 108 (1992), 1–10.
- [10] M. Kassi, A Liouville theorem for F-harmonic maps with finite F-energy, Electron. J. Differential Equations, 2006 (2006), 1–9.
- [11] Y. L. Xin, Geometry of harmonic maps, Birkhäuser, Boston, 1996.

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